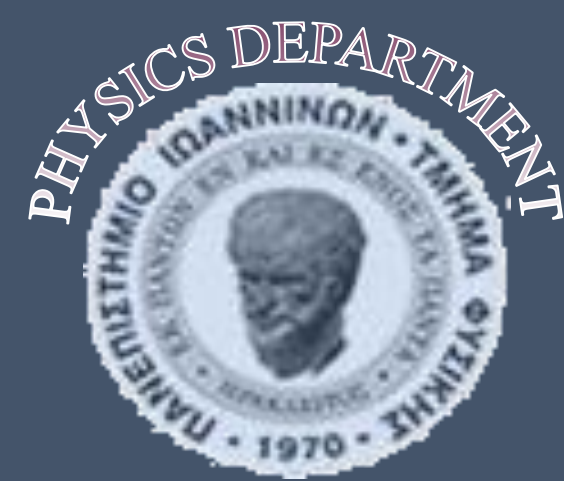


Identification of water fragmentation channels utilizing phase shifts under $\omega/2\omega$ fs laser field irradiation



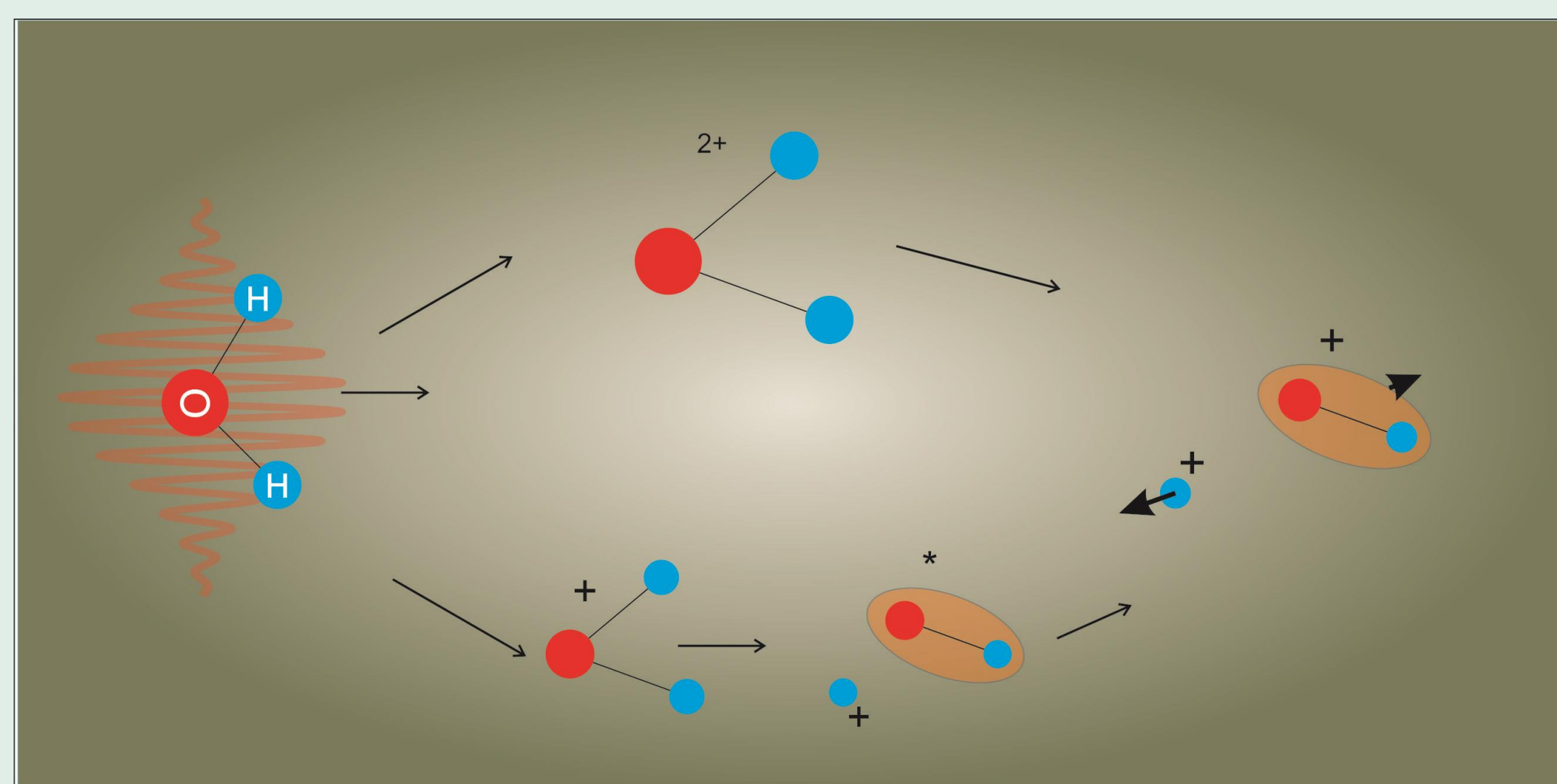
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Study and control of unimolecular reactions is feasible by laser irradiation using two frequencies ($\omega/2\omega$) [1].

- Ionization of molecules using strong $\omega/2\omega$ asymmetric fields can among other things:
1. Alter the branching ratio between two dissociative channels [2]
 2. Shed light on ultrafast phenomena like hydrogen migration [2,3]
 3. Distinguish between dissociative states that release similar fragments [4].

In the present work, we have studied the two-body fragmentation H_2O^{2+} according to the reaction:



Two body fragmentation routes of water leading to $H^+ (E_{kin} = 4.8 \text{ eV}) + OH^+ (E_{kin} = 0.3 \text{ eV})$ (a)

- Water dications' fragmentation was investigated at two $\omega/2\omega$ wavelength regions 1250/625nm and 1400/700nm, generated by an Optical Parametric Amplifier (seeded by a 20 fs Ti: Sapphire laser) and an additional BBO crystal.
- The phase between the ω and 2ω components is controlled by rotating a Calcite plate in the path of the second harmonic in steps of 0.02° .
- The ionization signal is recorded by means of time of flight mass spectrometry.

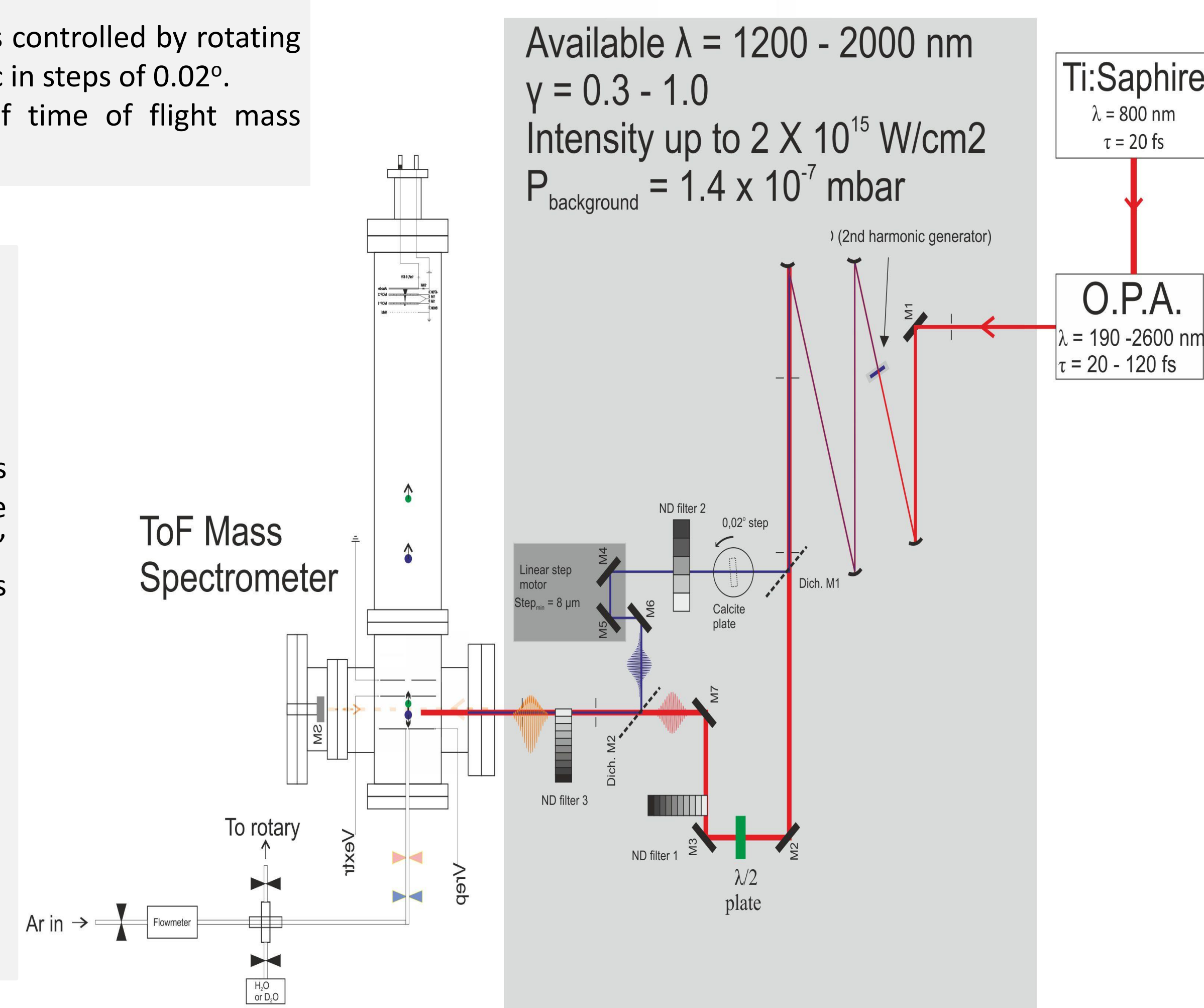
The composed field is mathematically described as:

$$E_{tot}(t, \phi) = E_0(t)[\cos(\omega t) + \gamma \cos(2\omega t + \phi)],$$

where γ is the ratio of the 2ω and ω electric fields and ϕ is the phase between the two beams. One useful parameter that is indicative of the fragments' ejection direction is the asymmetry parameter and is defined as

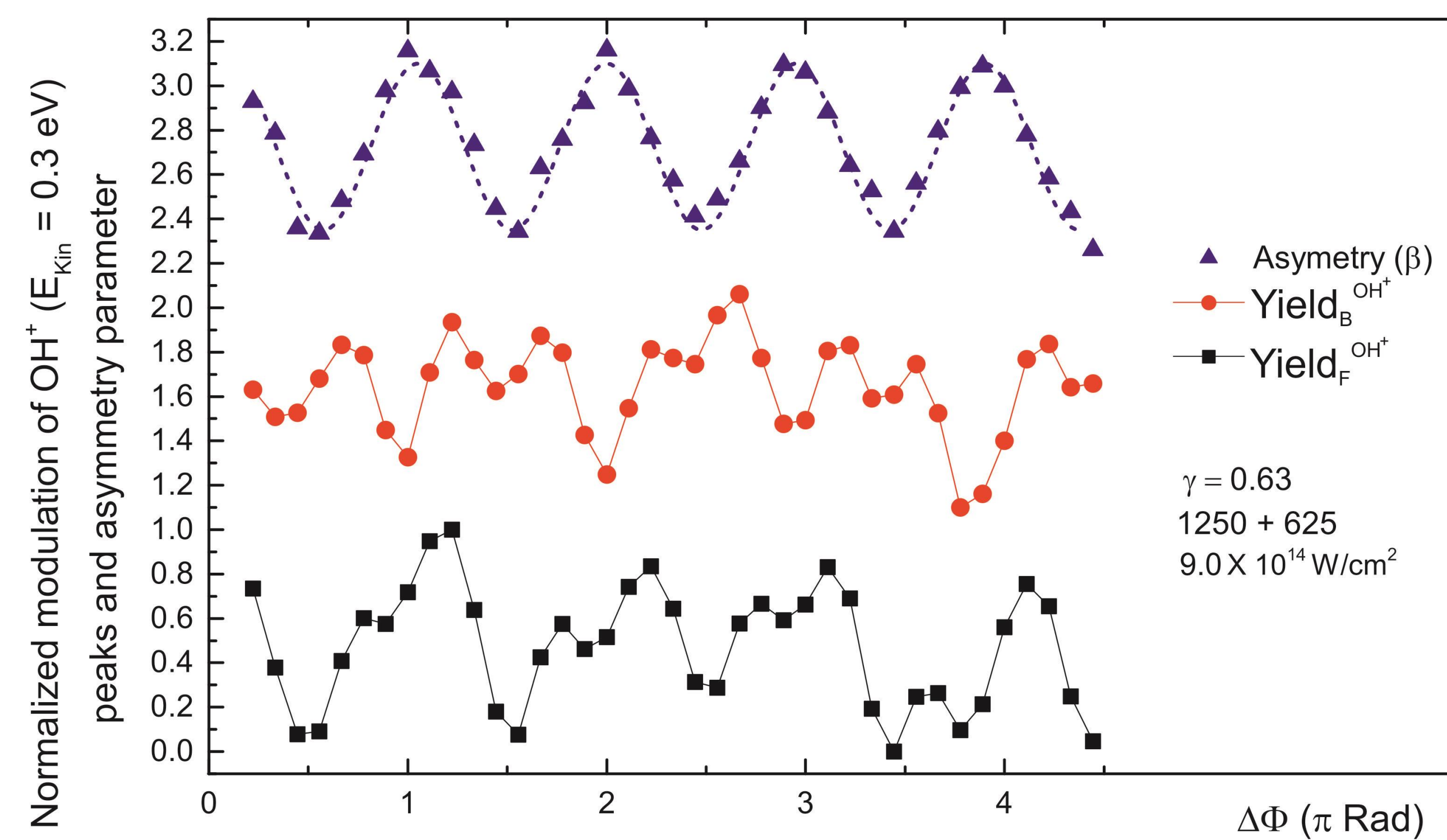
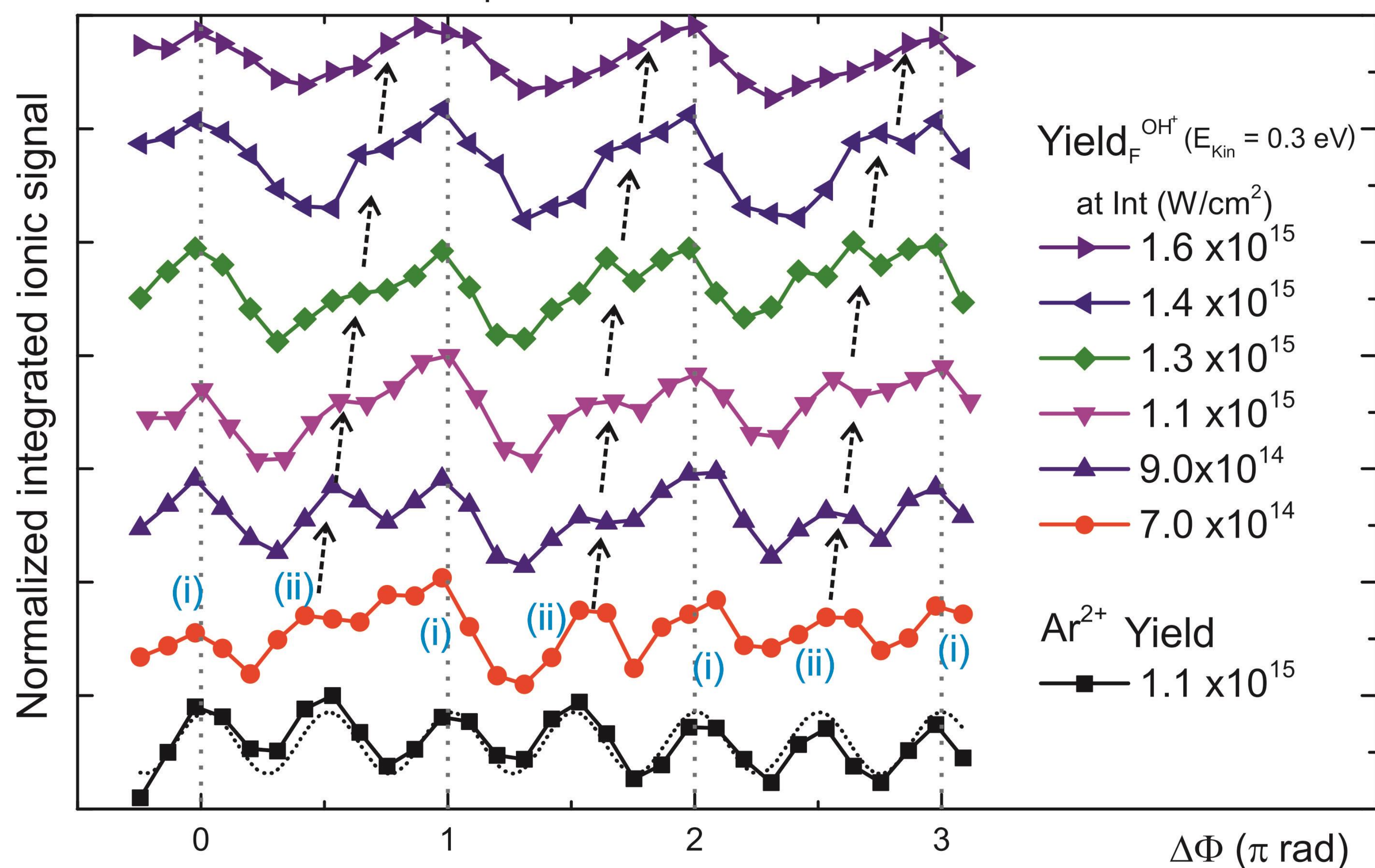
$$\beta = \frac{Yield_{Forward} - Yield_{Backward}}{Yield_{Forward} + Yield_{Backward}}$$

Experimental Setup ($\omega/2\omega$)



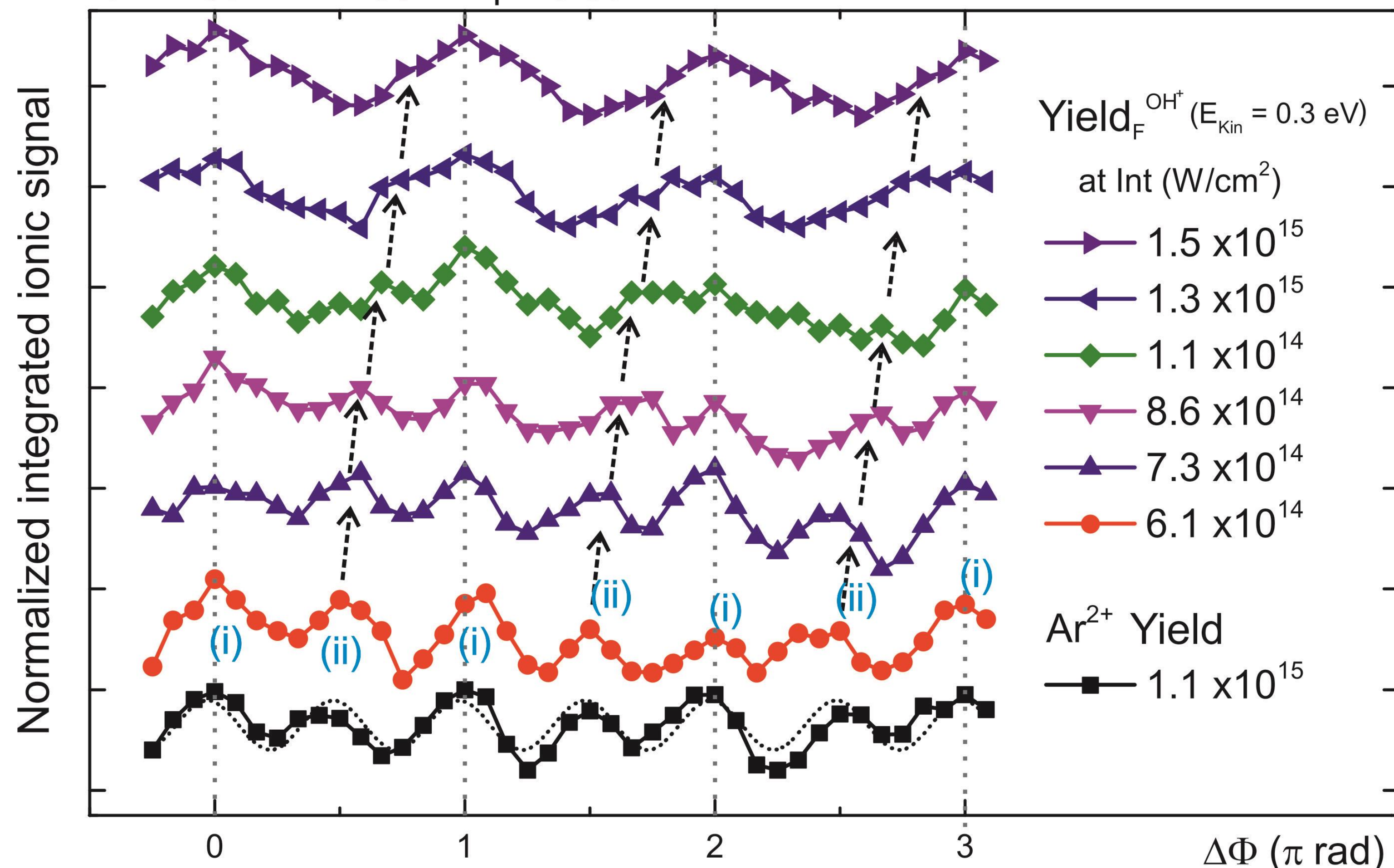
- A phase difference $\phi = \pi$ between $Y_F^{OH^+}$ and $Y_B^{OH^+}$ components was observed while the phase difference with respect to the components of the $H^+ (E_{kin} = 4.8 \text{ eV})$ peak was also $\phi = \pi$ [2]. This observation, in conjunction with momentum conservation confirms that these fragments are released from H_2O^{2+} .
- The asymmetry parameter (β) of the OH^+ peak exhibits a sinusoidal modulation that is dependent on the asymmetric field's direction (ω behavior).
- The individual OH^+ peak components ($Y_F^{OH^+}$ and $Y_B^{OH^+}$) show a dependence on the double frequency (2ω behavior).

$\lambda = 1250 \text{ nm} + 625 \text{ nm} \quad \gamma = 0.63$



- Ar^{2+} peak is monitored for phase calibration
- The phase difference between the OH^+ peak components ((i) and (ii)) is dependent of the laser's intensity and is perceived like a phase shift as the intensity rises
- This shows that two precursor dissociative states lead to these OH^+ fragments, despite the fact that they are ejected with the same (within the experimental error) kinetic energy.

$\lambda = 1400 \text{ nm} + 700 \text{ nm} \quad \gamma = 0.63$



- Two dissociation channels release similar fragments. The first channel takes place in a dication's state while the second one via an energetically lower lying autoionizing state of H_2O^+ , which after deprotonation, fragments to OH^* leading finally to OH^+ [5].
- Direct ionization (non-sequential) is favored at relatively lower laser intensities (peak (i)), while the contribution of the sequential (rescattering) process is increased with laser intensity (peak (ii)).
- For the higher laser intensities applied, the phase difference between the individual OH^+ peak components is small to be distinguishable and the ω behavior dominates.

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