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Multiscale analysis of chemical vapor deposition of copper

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1 Introduction

Chemical Vapor Deposition CVD is a widely used process for producing thin solid films and, nanoscale structures from gaseous reactants. Its applications range from coatings for wear and corrosion resistance, high temperature protection and erosion protection, to semiconductors and related devices (integrated circuits, sensors and optoelectronic devices), optical fibers for telecommunications, composites, powders and nanomachines, nanostructures, and nanowires[1]. Compared to physical deposition techniques (PVD), such as sputtering, sublimation, evaporation and other chemical based deposition techniques such as atomic layer deposition (ALD), CVD has the ability to produce a large variety of films and coatings of metals, semiconductors, and inorganic as well as organic compounds, possessing reproducible and controllable properties. The latter include film composition, uniformity, surface morphology, purity and uniform thickness on highly irregularly shaped surfaces. These properties vary with the application and material and they strongly depend on the usage of the fabricated device.

CVD is performed in specially designed reactors, the CVD reactors, equipped with the wafers or substrates, where the film is deposited through surface reactions. The gaseous reactants, including the precursor i.e., the compound that carries the chemical species to be deposited, enter the CVD reactor and are guided by gases, the carrier gas, to the wafer where the surface reactions and the deposition of the film occurs.

CVD is a multiscale process: From the macro (or bulk or gas phase) of the CVD reactor – of the order of some cm or even m – to the micro-scale of a predefined topography on the wafer's surface – of the order of μm - or the nano-scale of the growing surface of the film which is of the order of nm or \AA different processes are performed. Depending on the level of detail, dictated by the underlying theoretical or technological question, someone can study a CVD process in one of these scales or perform multiscale analysis.

The present study focuses on the multiscale modeling of Copper (Cu) CVD from a novel precursor, namely from copper(I) N,N'-di-isopropylacetamidinate $[\text{Cu}(\text{iPr}-\text{Me}-\text{amd})_2]$ or $[\text{Cu}(\text{amd})_2]$ where $\text{amd} = \text{CH}(\text{CH}_3)_2\text{NC}(\text{CH}_3)\text{NCH}(\text{CH}_3)_2$ [2]. Cu tends to replace Al in the interconnection of modern integrated circuits because of improved properties against Al. Specifically, Cu shows lower resistance, twice the thermal conductivity, 100 times greater resistance to electromigration which eventually can lead to material failure[3]–[5]. The deposition of Cu is yet challenging, especially in industrial scale since robust methods for depositing Cu films

are still unknown or the existing methods are insufficient due to lack of appropriate materials. $[\text{Cu}(\text{amd})]_2$ was proposed by [2] as a Cu precursor because of its properties against other Cu precursors such as low deposition temperatures, it is stable in the presence of other substances containing oxygen or halogens and high thickness uniformity [6].

In our previous work [6], based on the experimental work of [7], we proposed a deposition mechanism depending on the thermal region of deposition. In low temperature regions ($T < 506$ K) an Arrhenius type kinetics was proposed and for higher temperatures ($T > 506$ K) a Langmuir–Hinshelwood (LH) mechanism. The different reaction kinetics were proposed to capture the decrement of the deposition rate as the temperature increased. Similar behaviors have been observed previously for other systems, e.g. Al systems, but this decrement was attributed to the existence of volumetric reactions – i.e. reactions that perform in the bulk phase of the CVD reactor and not in the substrate – which result in the depletion of the precursor. The experimental work of [7], suggested that in the case of $[\text{Cu}(\text{amd})]_2$ there is only a single surface reaction taking place in the substrate [see Eq. (S2.1)] and no volumetric reactions up until 573 K. For that, based on previous works [8], [9], for the deposition of Cu from another precursor, namely $\text{Cu}(\text{hfac})_2$, we proposed a LH reaction mechanism which takes into account inhibition effects from amd. Very recently, Spencer et al. [10] proposed a detailed model which assumes volumetric reactions in the entire temperature deposition spectrum. Their computational results are in an excellent agreement with the experimental data. Nevertheless, their work did not cover multiple scales and for that could not explicitly take into account amd inhibition effects.

In this work we explore if amd can act as an inhibitor and reduce the deposition rate in high temperatures in a CVD multiscale context. In order to explore the inhibition effects by amd we link the computations in the macro-scale of a CVD reactor with a kinetic Monte Carlo model which takes into account explicitly the inhibition effects by amd. The macro-scale computations provide the mass fluxes of the different reactants reaching the substrate surface which are then used by the KMC software to perform the various processes that mimic the growth of Cu. Furthermore, we performed nonlinear computations in the macro-scale model of the CVD reactor in order to investigate if the transport phenomena can affect the deposition rate at high temperatures and we expand our previous nonlinear computations [11] to investigate the effect of the self-sustained periodic solutions to the deposition rate and films thickness non-uniformity.

2. Modelling the macro-scale of the CVD reactor

Due to the importance of the CVD processes in academia and industry, macro-scale CVD reactor models are well established through a series of work during the past four decades[12]–[18]. These models are based on the conservation laws of mass, momentum, energy, and species mass fraction and are described by partial differential equations[1].

2.1 Governing equations

At steady state and in vector notation the aforementioned set of equations is,

a) Continuity equation

$$\nabla \cdot (\rho \mathbf{u}) = 0 \quad (2.1)$$

where ρ is the density of the gas mixture and \mathbf{u} the velocity.

b) Momentum equation

$$\begin{aligned} \nabla \cdot (\rho \mathbf{u} \mathbf{u}) = & -\nabla P + \nabla \\ & \cdot \left[\mu (\nabla \mathbf{u} + \nabla \mathbf{u}^T) \right. \\ & \left. - \mu \frac{2}{3} (\nabla \cdot \mathbf{u}) \mathbf{I} \right] + \rho \mathbf{g} \end{aligned} \quad (2.2)$$

where P is the pressure, μ the viscosity, \mathbf{I} the unit tensor, and \mathbf{g} is the gravity acceleration.

c) Energy equation

$$\begin{aligned} C_p \nabla \cdot (\rho \mathbf{u} T) = & \nabla \cdot (\lambda \nabla T) \\ & - \sum_{i=1}^{N_{species}} \mathbf{j}_i \cdot \frac{\nabla H_i}{M_i} \\ & - \sum_{i=1}^{N_{species}} \sum_{k=1}^{N_r} H_i \gamma_{ik} r_k^g \end{aligned} \quad (2.3)$$

where C_p is the specific heat of the gas mixture, T the temperature, λ the thermal conductivity, \mathbf{j}_i the diffusive mass flux of species i , H_i the enthalpy of formation of species i , M_i the molecular

weight of species i , γ_{ik} the stoichiometric coefficient of species i in reaction k , $N_{species}$ the number of species in the gas phase, N_r the number of volumetric reactions, and r_k^g is the molar rate of volumetric reaction k .

d) Species transport equation

$$\begin{aligned} \nabla \cdot (\rho \mathbf{u} y_i) &= -\nabla \cdot \mathbf{j}_i + M_i \sum_{k=1}^{N_r} \gamma_{ik} r_k^g, i \\ &= 1, \dots, N - 1 \end{aligned} \quad (2.4)$$

where y_i is the mass fraction of species i . The aforementioned system of equations closes with the ideal gas law. The latter equations are discretized and solved with both CFD software Fluent [19] and Comsol [20]. Fluent model is used for the multiscale modeling while the Comsol model for the nonlinear analysis (see details in Section 5).

2.2 Reactor geometry, species mixture properties & boundary conditions

The macro-scale model is a 3D stagnation flow, cold wall, CVD reactor described in detail in previous works [6], [10], [21], [22] (see Fig. 2.1a,b). A 3D model is necessary to describe the phenomena inside the reactor because of the showerhead, which is used to homogenize the flow above the substrate, consists of 1450 holes of 0.76 mm diameter (see Fig. 2.1b and d). The thickness of the plate is 3 mm. A mesh of 1.232.468 cells comes out to be an appropriate one based on mesh independent study. Due to the size of the resulting computational problem, parallel processing in computer clusters is indispensable. The CPU time and memory required for the solution of 3D problem are ~4 h and 2 GB, respectively in Fluent software.

For the case of Cu CVD, the properties of the individual species and of the mixture are computed as in [6]. The Lennard-Jones (LJ) parameters, namely σ and ε , are the parameters of the LJ potential and are needed for the estimation of the properties in the gas phase of the CVD reactor. σ is the measure of the size of the molecules and ε/k is a measure of how strongly the molecules attract each other. For the unknown species, namely $[\text{Cu(amd)}]_2$ and H(amd) , their

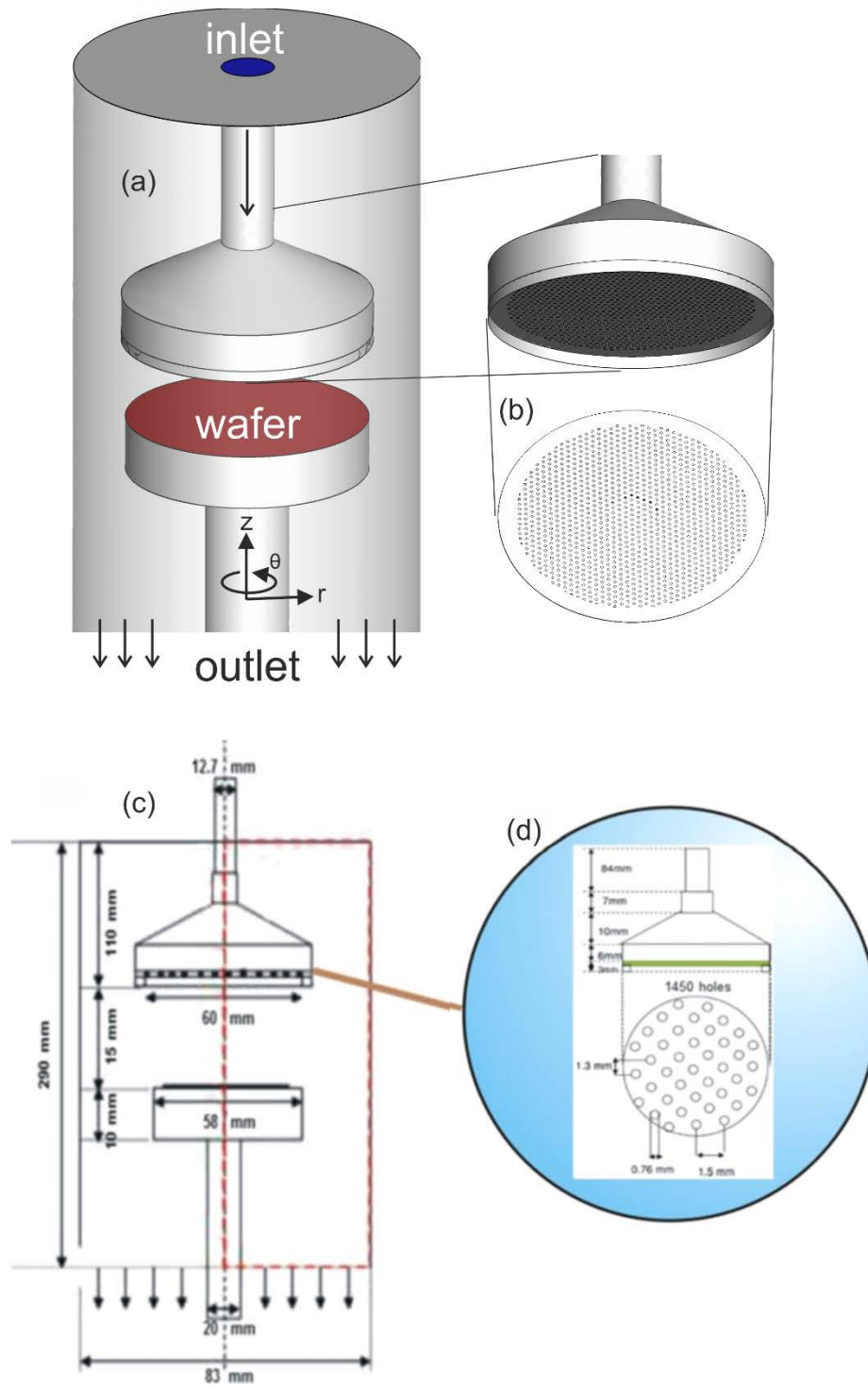


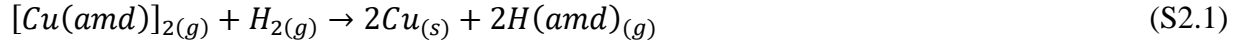
Figure 2.1 (a) The bulk phase of the CVD reactor (b) Details of the showerhead which is used to homogenize the flow above the wafer (c) Dimensions of the CVD reactor and (d) dimensions of showerhead.

values are calculated with group contribution methods and for $[\text{Cu}(\text{amd})]_2$ are: $\sigma = 10.8525 \text{ \AA}$ and $\varepsilon/k = 423.2 \text{ K}$ and for $\text{H}(\text{amd})$: $\sigma = 9.4874 \text{ \AA}$ and $\varepsilon/k = 534.8 \text{ K}$.

Concerning the boundary conditions, for the velocity: at the inlet of the reactor a constant mass inflow rate of $7.473 \times 10^{-6} \text{ kg/s}$ is imposed. No-slip condition is imposed at all the walls of the reactor. At the outlet, a standard outflow boundary condition is used. For the species: the mass fractions of the species entering the reactor are $y_{[\text{Cu}(\text{amd})]_2} = 0.001016$, $y_{\text{H}_2} = 0.004107$ and $y_{\text{N}_2} = 0.2556$; the rest is Ar. The flux of all species at the reactor walls is zero with an exception on the substrate where the surface reactions (deposition) take place. An overall mass balance correction is imposed at the outlet. The operating pressure of the reactor is 1333 Pa. The wafer temperature, T_w , varies as 473 K, 493 K, 513 K, 533 K, 553 K, 573 K, 593 K and 623 K, according to the available experimental data.

2.3 Surface reaction kinetics in the macro-scale model

Based on the experimental work of [7], we assume only a single surface reaction which leads to the deposition of Cu, namely,



Two reaction kinetics expressions are used; for $T_w = 473 \text{ K}$ to 513 K an Arrhenius expression is used and from 533 K to 623 K a Langmuir-Hinshelwood (LH) expression. The Arrhenius expression reads,

$$r_A = k_A \exp\left(-\frac{E_1}{RT}\right) C_{[\text{Cu}(\text{amd})]_2} C_{\text{H}_2} \quad (2.5)$$

where $k_A = 1.33 \cdot 10^{10} \text{ KJ}/(\text{m}^2 \text{ mol})$ and $E_1 = 66 \text{ KJ/mol}$ and the LH reads,

$$r_{LH} = \frac{k_a k_d C_{[\text{Cu}(\text{amd})]_2} C_{\text{H}_2}}{(2k_a C_{[\text{Cu}(\text{amd})]_2} + k_c C_{\text{H}(\text{amd})}) + k_d C_{\text{H}_2}} \quad (2.6)$$

Eq. 2.6 takes implicitly into account inhibition effects from H(amd). The parameters $k_d = r_A$, $k_a = 0.07 \left[\frac{m}{s} \right]$ and $k_c = 0.01 \left[\frac{m}{s} \right]$ are fitted to the experimental data. None of the above kinetics is valid on the entire spectrum of T_w ; the Arrhenius type is valid on low and the LH on high temperatures.

3. Monte Carlo & Kinetic Monte Carlo methods in deposition processes

Monte Carlo (MC) and kinetic Monte Carlo (kMC) are widely used methods in many fields of science and engineering: From materials science and polymers properties [23], astrophysics and black holes mergers [24] to computational geometry and volume approximation [25]. Their popularity in materials science stems from their inherent ability to simulate the molecular level of materials seamlessly. In MC/kMC, the particles (molecules, atoms, beads) move stochastically according to specific rules (events/processes), transferring the system randomly over the phase-space and approximating the mean values of various properties. In contrast to other molecular methods such as molecular dynamics (MD), the system in MC cannot easily be trapped in local energy minima and even if it is trapped, it can be “kicked out” to other states by incorporating sophisticated events. Furthermore, kMC filters out vibrational movements, allowing it to run over much larger spatial and time scales than MD. Especially in film growth, snapshots of MC/kMC simulation can be directly related and compared to scanning tunneling microscopy images. MC/kMC have notable applications in the study of film deposition processes—probably the most important in the fabrication of semiconductor devices.

During any deposition process the material is deposited upon a surface either by physical or chemical processes. In physical processes, the material is injected as a gas and sticks (adsorbs) on the deposition surface with a probability (or rate). PVD, which encompasses sputtering and evaporation, belongs to this category of deposition processes. In chemical processes, the material is grown on the deposition surface through surface reactions.

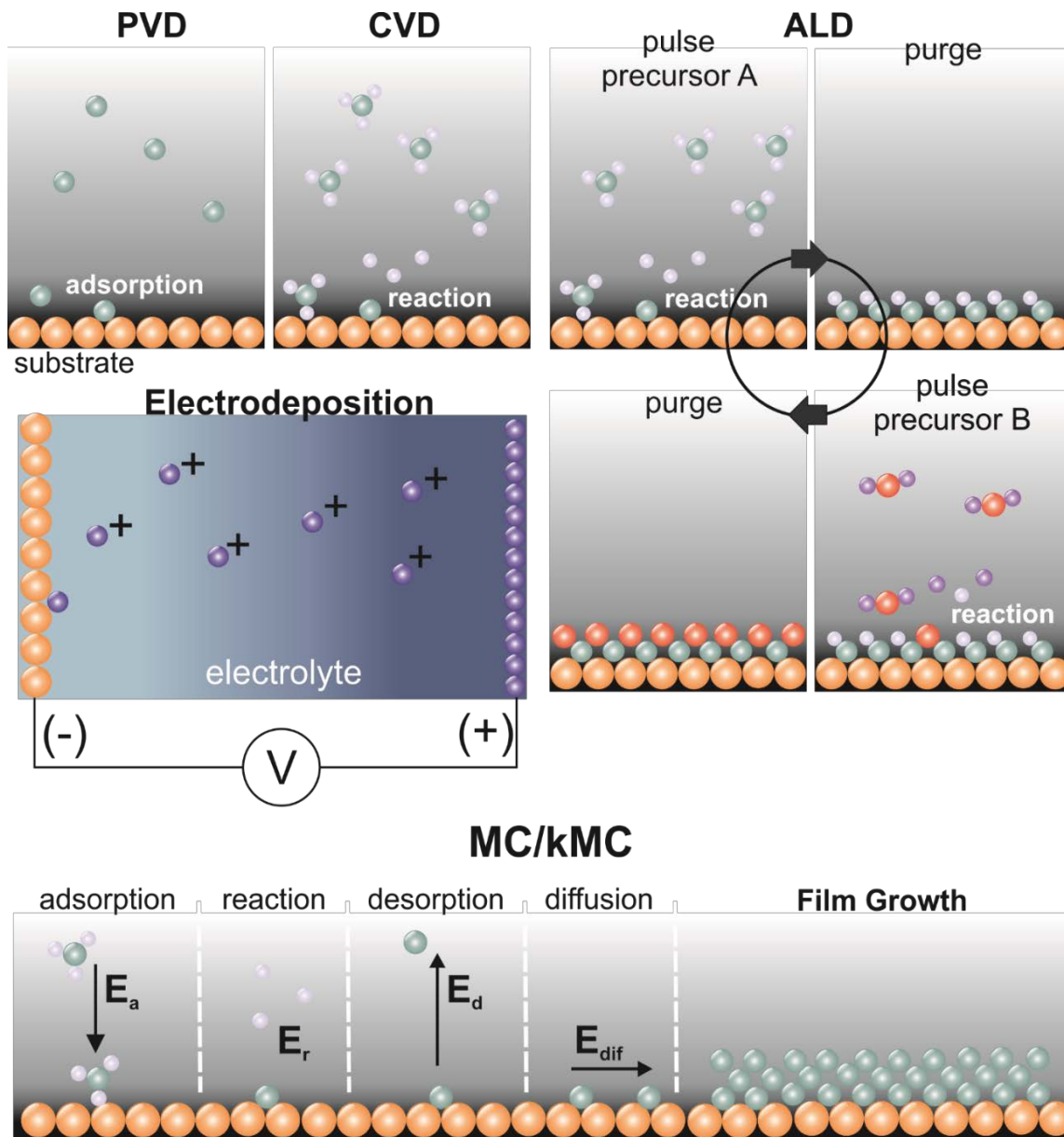


Figure 3.1 Schematic representation at the molecular level of the basic principles in deposition processes. In PVD, particles are adsorbed in the substrate. In CVD, particles (here an arbitrary molecule is shown) react on the substrate surface to grow the film. In ALD, multiple precursors (here A & B) are injected into the reactor though pulses in a cyclic way of pulses-purges and grow the film through self-limiting surface reactions. In electrodeposition, the voltage (V) is applied causing the particles from the cathode to move to the anode through an electrolyte where the film is grown through adsorption or surface reactions. In MC/kMC methods, an activation energy is assigned to each event. In MC, if the event leads the system to a smaller energy, it is accepted unconditionally. If not, it is selected with a probability according to Eq. 3.2. In kMC, an event is selected from a predefined rate catalog containing all possible rates and performed unconditionally. Then the system evolves in time according to Eq.3.8. The four basic events – adsorption, reaction, desorption and diffusion- whose combination can describe a deposition process, with their activation energies, E_a , E_r , E_d and E_{dif} are shown.

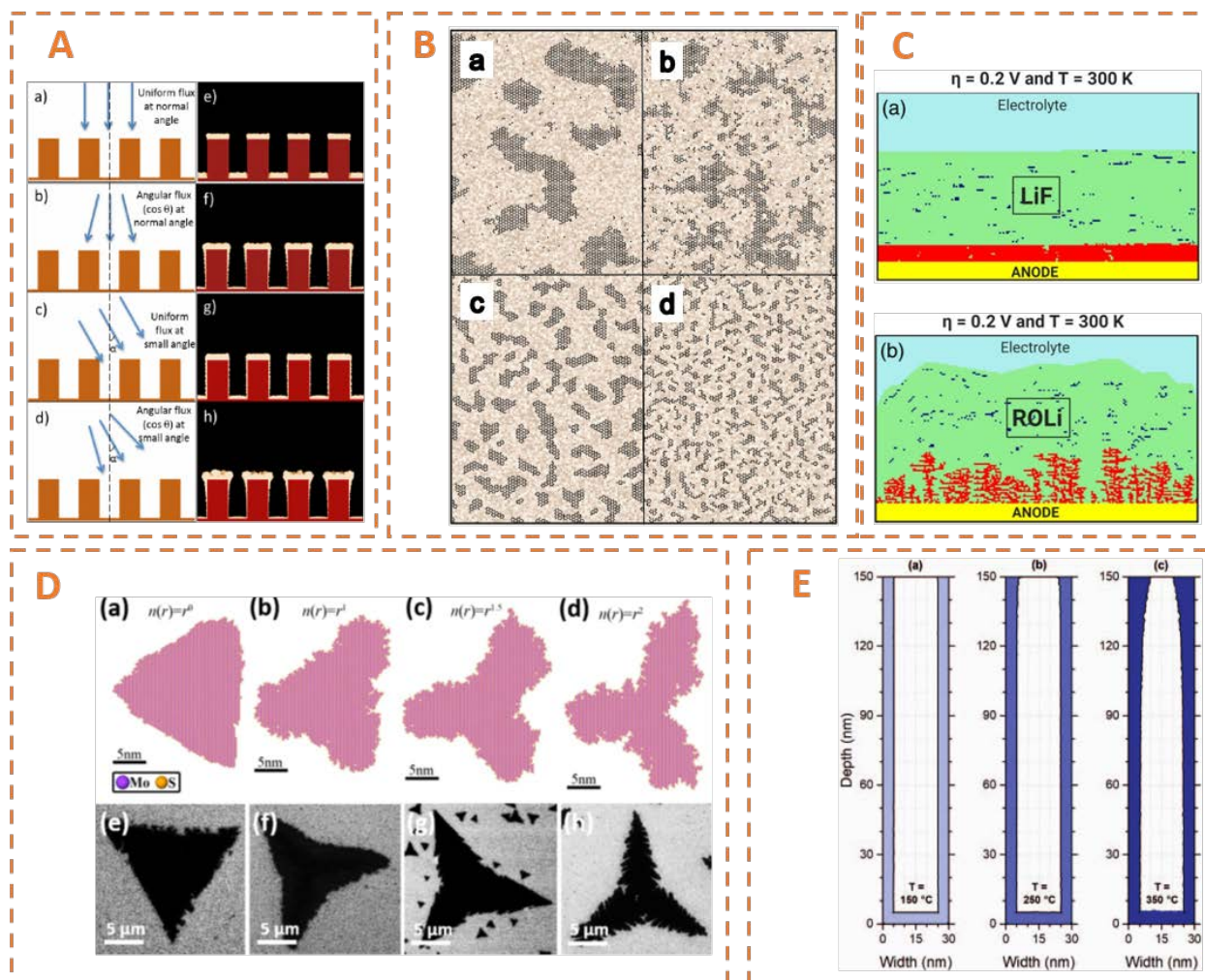


Figure 3.2 **A.** PVD of Ag core/shell on In_2S_3 nanorods: MC results for four different conditions: (a & e) Uniform incident flux, $\alpha = 0^\circ$ (α is the angle between the incoming flux and the surface normal to the substrate), (b & f) angular flux $\sim \cos(\theta)$ with $\alpha = 0^\circ$, (c & g) uniform flux with $\alpha = 30^\circ$ and (d & h) angular flux $\sim \cos(\theta)$ with $\alpha = 30^\circ$ (reprinted with permission from ref. [26]) **B.** Graphene CVD: MC results for varying strength of surface roughness parameter ζ and mobility parameter T_s for (a) $\zeta = 1.2$, $T_s = \infty$ (b) $\zeta = 1.5$, $T_s = \infty$ (c), $\zeta = 1.5$, $T_s = 0$ and (d) $\zeta = 1.2$, $T_s = 0$. The color represents roughness (reprinted with permission from ref. [27]). **C.** Electrodeposition in Li metal anode: kMC results for (a) uniform Li deposition in LiF electrolyte, (b) dendrite growth in ROLi electrolyte. The red particles represent the Li metal which is being deposited, and the dark-blue particles represent the Li^+ ions which are diffusing. The process conditions (η - activation overpotential and T - temperature) are also shown (reprinted with permission from ref. [28]). **D.** MoS_2 CVD: kMC results for the morphology of MoS_2 as a function of the Mo density gradient, $n(r) - r$ distance from Mo source (a) no Mo density gradient, $n(r) = r^0$ (b-d) $n(r) = r^1$, $n(r) = r^{1.5}$, and $n(r) = r^2$ (e-h) corresponding SEM images (reprinted with permission from ref. [29]). **E.** SiN_x ALD: Thickness profile in a microstructure with aspect ratio equal to 5 for parasitic CVD (no-ALD) reactions at different T_w (a) 150°C , (b) 250°C and (c) 350°C , determined by kMC computations (reprinted with permission from ref. [30]).

Examples of chemical deposition processes are CVD, ALD and electrochemical deposition. In the context of MC/kMC methods the different processes are shown schematically in Fig. 3.1.

Concerning the recent works of MC and kMC on deposition processes, most PVD computational works deal with the growth of nanorods[26], [27], [31]–[33] while in CVD, most works examine the growth of graphene[29], [34]–[40], followed by 2D materials and namely Transition metal dichalcogenides (TMDs - MoS₂, WSe₂ and WS₂)[41]–[46]. kMC models have also been developed for the deposition of diamond [47], GaAs [48] and AlN [49]film, plasma enhanced a-Si:H CVD [50], hybrid MD/kMC [51] and for growth in extreme pressure conditions [52]. Concerning ALD, most works deal with the study of step coverage inside micro-/nano-features,[30], [53]–[56]. In the context of electrodeposition, most work perform MC/kMC to study the properties of Li anodes metal batteries[57], [58]. Representative results from the MC/kMC models applied in recent works for deposition processes are shown in Fig. 3.2.

3.1 The kMC method

In MC and kMC algorithms, sequential events are performed stochastically. MC solves the steady state Master Equation (ME) and kMC the transient one. The transient ME reads,

$$\frac{\partial p_j(t)}{\partial t} = \sum_{i \neq j} p_i(t) T_{ij} - \sum_{i \neq j} p_j(t) T_{ji} \quad (3.1)$$

where $p_{j(i)}$ is the probability of the system to be found in state j (i) at time t . T_{ij} and T_{ji} denote the transition rate or transition probability from state i to j and vice versa. Each event occurs at a certain probability/rate to form a Markov chain[59]. To generate the Markov chain the desired probability distribution $p_{i(j)}$ must obey the detailed balanced condition,

$$p_i T_{ij} = p_j T_{ji} \quad (3.2)$$

$p_{i(j)}$ are Boltzmann distributions,

$$p_{i(j)} = \exp\left(-\frac{E_{i(j)}}{k_b T}\right) \quad (3.3)$$

where $E_{i(j)}$ is the energy of the system in state $i(j)$. [60] proposed that,

$$T_{ij} = \begin{cases} 1, & \text{if } \Delta E \leq 0 \\ \exp\left(-\frac{\Delta E}{k_b T}\right), & \text{if } \Delta E > 0 \end{cases} \quad (3.4)$$

so that the system will unconditionally move from state i to j if $\Delta E < 0$ and with a probability $\exp\left(-\frac{\Delta E}{k_b T}\right)$ if $\Delta E > 0$, where $\Delta E = E_j - E_i$. Practically, a random number ξ is chosen between (0,1] and if $\xi < \exp(-\Delta E/k_b T)$, the system moves to the state j , otherwise the move is rejected. In this way, different states of the system are generated and the thermodynamic average of a quantity q_i reads,

$$\langle q \rangle = \frac{\sum_i q_i p_i}{\sum_i p_i} \quad (3.5)$$

KMC method solves Eq. 1. The most popular algorithm proposed by Bortz et al. [61] is termed as the N -fold method. In the N -fold method random transitions from i to j are performed unconditionally based on the transition rates, so that more likely transitions are selected more often. Every transition event i is assigned a rate which reads,

$$r_i = v_i \exp\left(-\frac{E_i}{k_b T}\right) \quad (3.6)$$

where v_i is a frequency prefactor, E_i is the energy barrier and T is the temperature. Practically, the simulation starts by defining all rates (rate catalog) r_i of all possible processes that describe the physical problem. The total rate, $R = \sum_i r_i$, is first computed and then a process n is randomly chosen according to,

$$\sum_{i=1}^n \frac{r_i}{R} < \xi_1 < \sum_{i=1}^{n+1} \frac{r_i}{R} \quad (3.7)$$

where ξ_1 is randomly chosen in (0,1) and a single event is performed. The time advances as $t = t + \Delta t$ with Δt being,

$$\Delta t = -\frac{\ln \xi_2}{R} \quad (3.8)$$

where ξ_2 is an additional random number chosen in (0,1). R is recalculated based on the new system state. The algorithm stops when the desired time interval is reached.

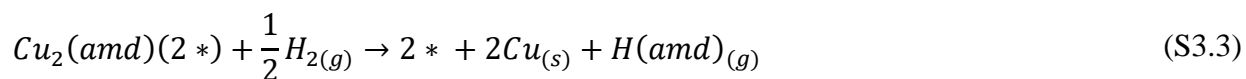
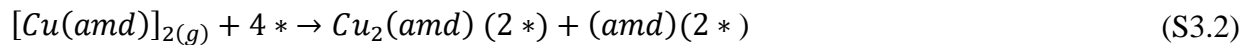
kMC rates must obey also the detailed balance condition (Eq. 3.2), even if the system is not in equilibrium, to ensure the dynamic evolution will correspond to a physical system [62]. kMC rates depend on both the particle and the lattice type that participate in the process and can be calculated via Transition State and Harmonic Transition State Theories (TST – HTST [62]), Density Functional Theory (DFT) and *ab initio* methods (e.g. [39], [41], [43], [44]).

Focusing on the deposition processes, the notion of lattice is of great importance. The lattice represents the deposition surface and is composed of sites upon which all events occur, simplifying the construction of the rate catalog. Depending on how the lattice is represented, the atomistic information can either be presented in full detail (e.g. [63]) or in a coarse-grained way where microscopic neighboring sites are coalesced into coarse cells (e.g. [64], [65]). Off-lattice kMC [66] methods have also been proposed in atomistic representation where the rate catalog is computed “on the fly” in every step.

3.2 kMC model of Cu deposition from [Cu(amd)]₂

The reaction mechanism proposed in this work is based on the first principles computations of [67]. According to their work, the most energetically favorable process is the dissociation of [Cu(amd)]₂ onto the substrate and the occupation of four positions of the Cu surface: two by Cu₂(amd) and two by amd (see Fig. 3.3). Thus, inhibitions effects are due to the adsorbed amd upon the surface. The H₂ above the substrate is in excess and according to our macro-scale computations, the mole fraction is ~1000 times more than [Cu(amd)]₂ (see Fig. 4.1), so we assume that H₂ does not limit the growth of Cu due to material inadequacy. As already stated, amd limits the growth rate because it blocks sites where [Cu(amd)]₂ can use to adsorb. The reaction of H₂

with amd leads to H(amd) which we assume that adsorbs instantly and releases potential sites for deposition. The latter mechanism is described as,



Here $Cu_2(amd)$ and amd are identical in terms of the reaction with H_2 since only the amd part reacts with H_2 .

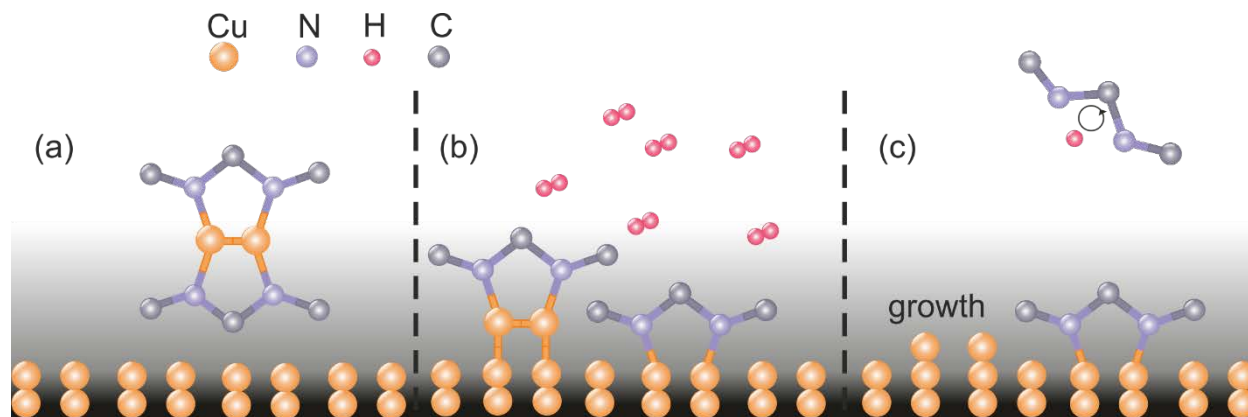


Figure 3.3 Schematic representation of the proposed reaction mechanism leading to the growth of Cu from $[Cu(amd)]_2$. (a) The $[Cu(amd)]_2$ molecule above the Cu surface. H on the $[Cu(amd)]_2$ molecule is not shown for simplicity. (b) Dissociation of $[Cu(amd)]_2$ molecule and inhibition effect from amd. $[Cu(amd)]_2$ cannot dissociate to sites occupied by amd. (c) Reaction with H_2 , desorption of $H(amd)$ and growth. The other amd still occupies two sites, preventing $[Cu(amd)]_2$ from dissociating to these sites.

The kMC model is based on Solid-On-Solid (SOS) approximation in a simple cubic lattice where $[\text{Cu(amd)}]_2$ dissociates upon four sites. We assume that the dissociation is performed instantly upon $[\text{Cu(amd)}]_2$ reaching the surface and thus the rate of the dissociation is the mass flux of $[\text{Cu(amd)}]_2$ reaching the substrate from the gas phase of the CVD reactor which reads[68],

$$r_1 = \frac{s_{[\text{Cu(amd)}]_2} \langle f_{[\text{Cu(amd)}]_2} \rangle P_{op}}{N_a C_{tot} \sqrt{2\pi M_{[\text{Cu(amd)}]_2} kT}} \quad (3.9)$$

where $s_{[\text{Cu(amd)}]_2}$ and $\langle f_{[\text{Cu(amd)}]_2} \rangle$ are the sticking coefficient and average over the wafer mole fraction of $[\text{Cu(amd)}]_2$ respectively, P_{op} the operation pressure of the reactor, N_a Avogadro's number, $C_{tot} = 2.0 \cdot 10^{19} [\text{sites}/\text{m}^2]$, the sites concentration on the surface of Cu, $M_{[\text{Cu(amd)}]_2}$ the molecular weight of $[\text{Cu(amd)}]_2$ and k Boltzmann's constant.

For the reaction of the adsorbed amd with H_2 the rate is given by,

$$r_2 = k_2 \exp\left(-\frac{n E_2}{R}\right) \quad (3.10)$$

$k_2 = 1.3 \cdot 10^{13} [-]$, $E_2 = 8.965 \cdot 10^4 \left[\frac{\text{kJ}}{\text{mol}}\right]$ and n is the number of nearest neighbors which are occupied by Cu ($n = 2,3,4,5,6,7,8$). That said, Eq. (3.10) takes into account the difficulty of the produced H(amd) to desorb and release the two sites for the dissociation of $[\text{Cu(amd)}]_2$. The pre-exponential factor and energy in Eq. (3.10) where fitted to the experimental data at 553 K from a single run and it was found that in order to capture the experimental data it should be a barrierless (athermal) process[69], [70]. H_2 is known to participate in barrierless processes as e.g. in the case of Si [71]–[73]. Still, such a claim must be investigated further via first principles, e.g. Density Functional Theory (DFT), computations[71] to better understand the mechanism and validate it.

3.3 The Apothesis software

There are many open-source software, based on the kMC method available online ([74]–[76]). However, these codes are primary for analyzing catalytic processes and not deposition, e.g.

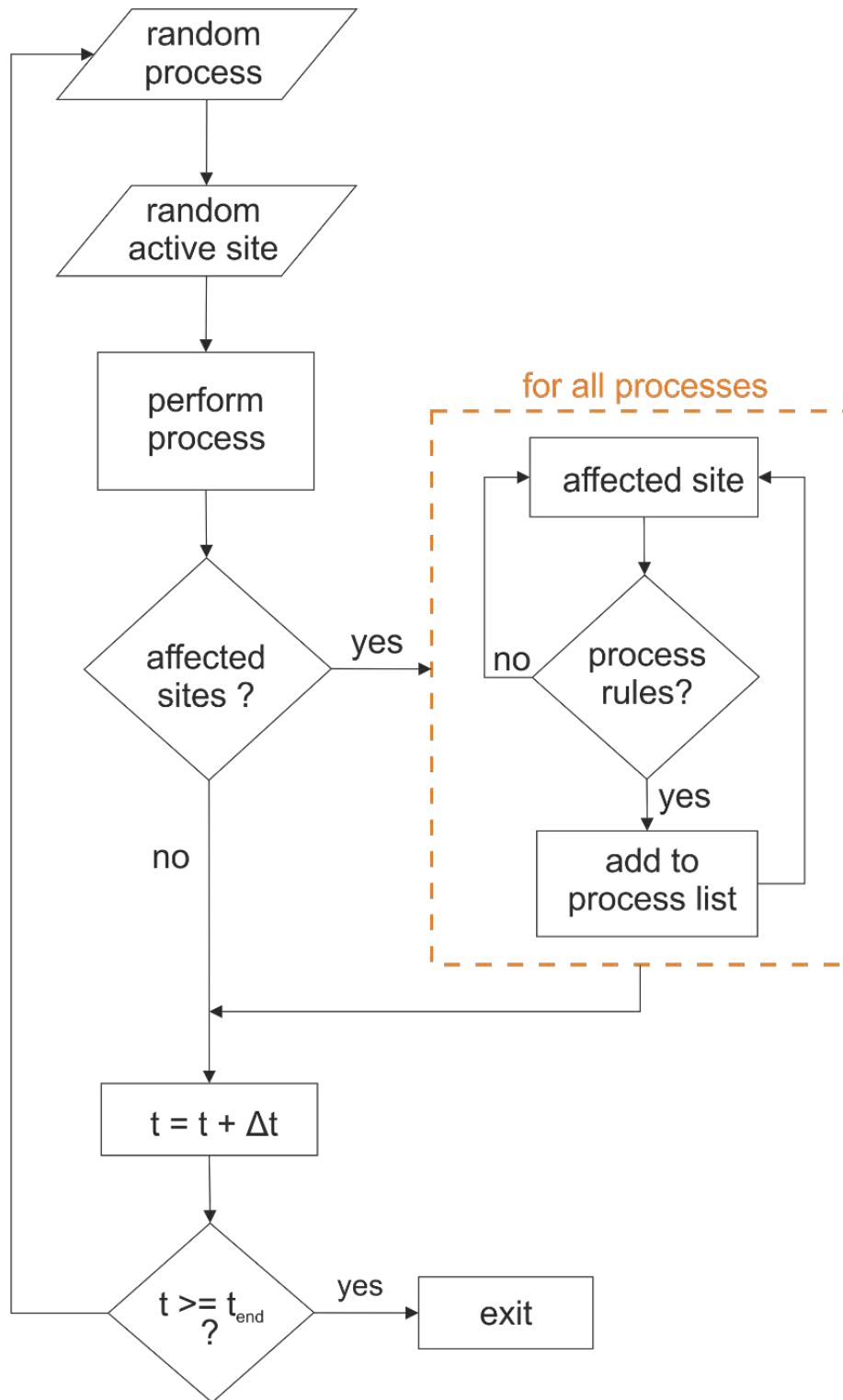


Figure 3.4 The computational “backbone” of Apothesis software. The algorithm starts by picking a random process and a random active site of this process. The process is performed and if it affects other sites a loop over all process is performed in order to update their lists according to their “rules” until all process lists are updated. After that the time is updated until the end time, t_{end} is reached.

they do not take into account the growth of the surface which is fundamental in deposition processes. Other software based on kMC[77], can be used only on the context of PVD since they cannot simulate chemical reactions. Other software with the same limitations have also been reviewed in [78].

In the context of this project, we developed the core of Apothesis. Apothesis is designed to be a generalized, pseudo-3D kMC software for deposition processes based on the solid-on-solid approximation. Apothesis is written in C++ and an expansion in python is under design which will enable users to add any process without re-compiling the code.

The main components of Apothesis are the lattice and the processes. The lattice is a predefined number of positions, termed as sites, where the different processes can perform on. Depending on the material it can be simple cubic, Body Cubic Centered (BCC), Face Cubic Centered (FCC) etc. and depending on the orientation (001), (011), (111) etc. A process is an act on a site of the lattice. If a process can perform on a site, this site is termed as “activated” otherwise it is “blocked”. An Apothesis process emulates a physical/chemical process e.g. the adoption of a species on an active site.

The sites of the lattice where a process can perform are kept on a list. The size of this list multiplied by the probability of the process is the process rate. In the course of a kMC simulation, each process list must be updated in order for the rates to be computed correctly. This makes this part of the code crucial in terms of the validity of the software and in terms of speed, since this is the most computationally demanding part of the kMC method in general. In order to generalize Apothesis, this part of the software should be very carefully designed and developed.

We tackled this issue by designing an abstract “Process” class which contains two primary functions, the “rules” and the “perform” functions. The “rules” function defines, computationally, when a process can be performed or not in a site and the “perform” is the process act on the lattice site. We take as an example the desorption of a particle of 4 nearest neighbors. The “rules” in this case would be to count the nearest neighbors of the site and if it is 4 then add it in its list. The “perform” would be to remove one particle (decrease site’s height) and mark the sites that are affected which in this case are all 4 nearest neighbors. A loop is performed to check, through the “rules”, if the affected sites must be removed or added to any of the other process lists. In that way, the screening is performed in only a small number of sites, and not the entire lattice, reducing significantly the computational time. In this way, we construct a generalized framework where

every process can be described independently of the lattice and thus, producing a robust computational environment for kMC computations. The computational “backbone” of Apothesis can be seen schematically in Fig. 3.4. Apothesis is freely available in github: <https://github.com/nixeimar/Apothesis> under GPL license.

4. Multiscale modeling of CVD: Coupling/Linking reactor scale with kMC models

The first description for the need of linking a reactor scale model of a CVD reactor with kMC model is found in the work of Srolovitz et al.[79] for the CVD of diamond. In the same year, Vlachos [80], presented a model that coupled the reactor scale model of a chemical reactor with a KMC model in catalytic processes and then extended it for CVD processes[81]. Raimondeau and Vlachos[82] presented a coupling of the scales in a vertical CVD reactor while at the same period Masi et al.[83] presented a multiscale model for the CVD of Si from SiH₄. They linked the models at different scales by assuming that the growth rate is independent of the scale that it is computed i.e. the deposition rate is computed from the reactor scale model and for this value of the deposition rate, the species mass fractions on the wafer surface are transferred to the KMC model for computing the roughness of the growing film. The same assumption was made by Grujicic and Lai [84] for performing multiscale computations for diamond growth.

In 2001, Lam and Vlachos [187] showed the effect of the operating conditions of a CVD reactor on the roughness development and the growth rate of the film. Dollet [85] developed an integrated multiscale framework for the CVD of silicon carbides. In a series of works[86]–[89], Cavallotti and Barbato et al. developed multiscale computational frameworks which was mainly applied to Si deposition. Masi et al.[90] used a multiscale–multi-hierarchy approach, based on a succession of simplified models, to describe the deposition and the film morphology evolution during the epitaxial silicon carbide deposition in an industrial hot wall reactor.

More recently, we linked the different scales of Si and Al CVD[91], [92], by assuming that the computed growth rate is independent of the scale that it is computed from, as in the work by Masi et al.[83]. Aviziotis et al.[93], based on the same framework, performed multiscale computations for the Fe CVD. Wang et al.[94] performed coupled simulations with an MC model to predict the columnar growth of Cu nanorods for different operating pressures of a plasma reactor

and Christofides *et al.*[28], [95]–[99] conducted multiscale simulations of plasma assisted ALD processes, combining CFD and kMC models to characterize and control the process.

4.1 Linking the scales

The linking of the scales in this work is performed in the boundary condition for the species equations. We use the term linking since we send only information from the macro-scale model to the kMC model. Since, the use of the reaction kinetics in the macro-scale (see Eqs. 2.5 & 2.6) produces results that are in fair agreement with the experimental data, it offers a realistic approximation of the values of $\langle f_{[Cu(amd)]_2} \rangle$. From $\langle f_{[Cu(amd)]_2} \rangle$, the mass flux of $[Cu(amd)]_2$, $F_{[Cu(amd)]_2}$, can be computed as,

$$F_{[Cu(amd)]_2} = \frac{\langle f_{[Cu(amd)]_2} \rangle P_{op}}{\sqrt{2\pi M_{[Cu(amd)]_2} kT}} \quad (4.1)$$

from which $s_{[Cu(amd)]_2}$ is estimated as [100],

$$s_{[Cu(amd)]_2} = \frac{\langle r_A \rangle (\langle r_{LH} \rangle)}{F_{[Cu(amd)]_2}} \quad (4.2)$$

where $\langle r_A \rangle$ ($\langle r_{LH} \rangle$) are the average values over the substrate. The values of $s_{[Cu(amd)]_2}$ and $f_{[Cu(amd)]_2}$, needed in Eq. (2.7) of the kMC code, are computed from the macro-scale model and serve as the linking point of the different scales. Since no volumetric reactions are assumed and a good approximation of the kinetics in the macro-scale exist, the present multiscale scheme avoids time consuming computations as in the case of coupling, i.e. the by-directional transfer of information between the scales, which would require an iterative scheme [85], [88], [90], [101].

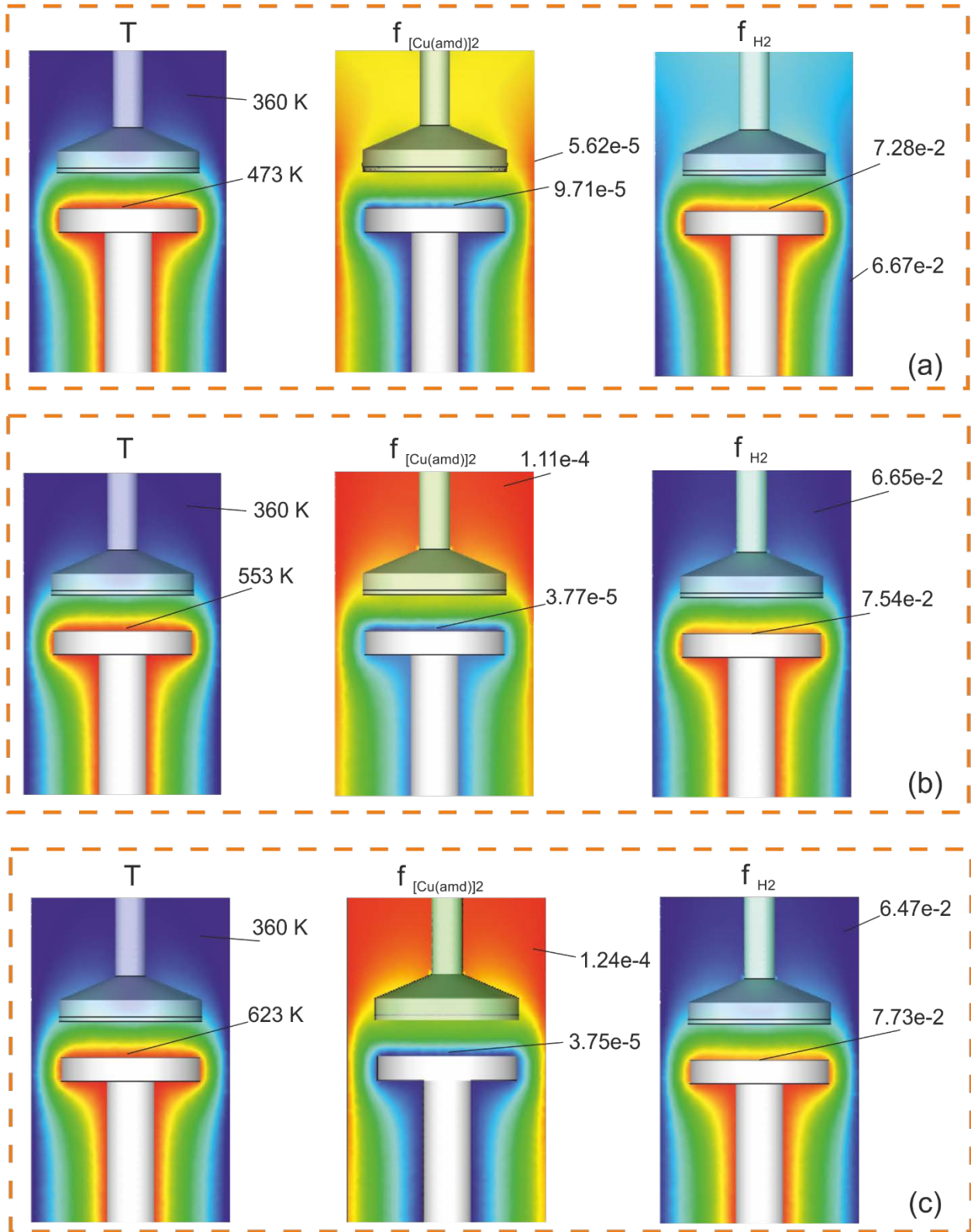


Figure 4.1 Temperature, $f_{[Cu(amd)]_2}$ and f_{H_2} distributions inside the CVD reactor for (a) $T_w = 473$ K, (b) $T_w = 553$ K and (c) $T_w = 623$ K. The maximum and minimum values are shown.

4.2 Results and discussion

Firstly, results from the computations in the macro-scale of the CVD reactor are presented. Three indicative T_w are selected based on the experimental data: 473 K where the lowest growth rate is observed, 553 K where the maximum growth rate is observed and 623 K where the maximum drop of the growth rate is observed as T_w increases. For these T_w the distributions of T , $f_{[Cu(amd)]_2}$ and f_{H_2} inside the reactor are shown in Fig. 3.1. The T distribution is almost identical for all T_w . Concerning the mole fractions distributions, $f_{[Cu(amd)]_2}$ in the cases of 533 K and 623 K is reduced compared to the 473 K case, which is expected since the reaction rate of $[Cu(amd)]_2$ is increased as T_w increases and more amount of $[Cu(amd)]_2$ is consumed. For the cases of 533 K and 623 K, $f_{[Cu(amd)]_2}$ is slightly reduced. Concerning f_{H_2} , there are no significant variations and it can be considered as constant. It is noted that it is $\sim 1,000$ times more than $[Cu(amd)]_2$ which validates our assumption that the amount of H_2 is not the limited step for the calculation of the growth rate in S2-S4.

The growth rate computed by our macro-scale computations and the two different reaction kinetics expressions vs T_w is shown in Fig. 4.2. As stated in our previous work [6], the Arrhenius expression prediction is in excellent agreement with the experimental data in low T_w , but as T_w increases it fails to predict the growth rate. On the other hand, the LH fails in low T_w , but satisfactorily predicts the experimental data in high T_w . Thus, none is valid in the entire T_w spectrum. It must be noted that both expressions are empirical. Nevertheless, since their combination can predict the growth rate in fair agreement with the experimental data, it can be used to compute the necessary information for the kMC model as states in Section 2.2.

The average value of $\langle f_{[Cu(amd)]_2} \rangle$ used in the computation of $s_{[Cu(amd)]_2}$ from the macro-scale computations are shown in Table 4.1 for every T_w . As T_w increases, up until T_s 513 K, $\langle f_{[Cu(amd)]_2} \rangle$ is decreased since the reaction rate increases. From 513 K until 573 K, $\langle f_{[Cu(amd)]_2} \rangle$ is practically constant and for 593 K and 623 K is slightly increased since the

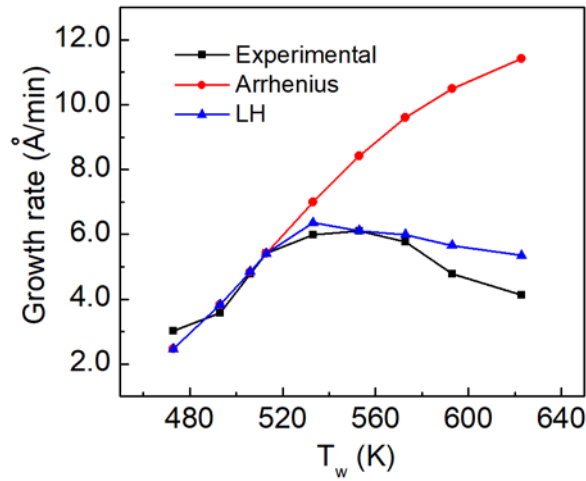


Figure 4.2 Deposition rate vs T_w as computed from the macro-scale model.

The two reaction kinetics, Arrhenius and LH are shown.

Table 4.1 The sticking coefficient $s_{[Cu(amd)]_2}$ and average, over the substrate area, mole fraction, $\langle f_{[Cu(amd)]_2} \rangle$ of $[Cu(amd)]_2$ computed in the macro-scale and used in the kMC model.

T_w [K]	473	493	506	513	533	553	573	593	623
$s_{[Cu(amd)]_2}$	0.027	0.048	0.069	0.099	0.097	0.096	0.094	0.093	0.091
$\langle f_{[Cu(amd)]_2} \rangle$ $\times 10^5$	5.94	5.12	4.55	3.96	3.98	3.99	3.99	4.01	4.02

growth rate (and the reaction rate for the particular case) is decreased. For T_w 593 K and 623 K, $\langle f_{[Cu(amd)]_2} \rangle$ is slightly increases but $s_{[Cu(amd)]_2}$ is decreased since in order to yield the decreased deposition rate in these T_w .

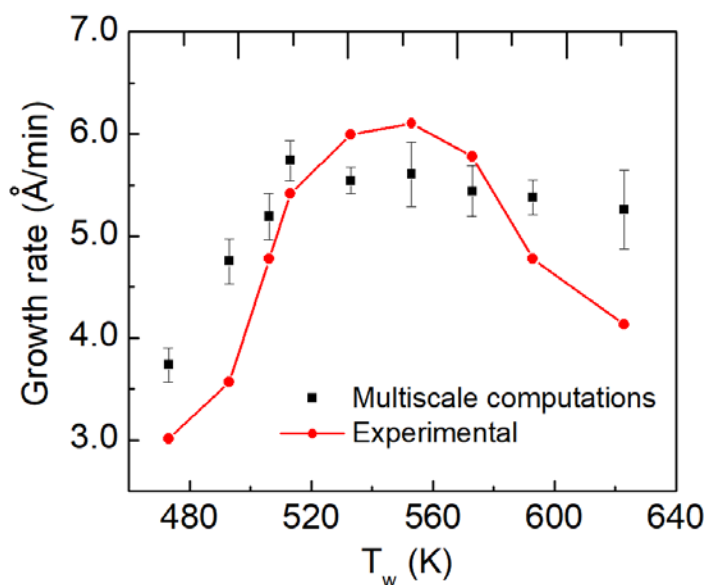


Figure 4.3 Growth rate vs T_w via multiscale computations. Information from the macro-scale of the CVD reactor is passed to the kMC model.

Using the data from Table 4.1 as input for the kMC model which explicitly takes into account inhibition effects from amd, described in Section 3.2, we predict the growth rates for every T_w shown in Fig. 4.3. In Fig. 4.3, the average growth rate over 10 kMC runs in a 300x300 lattice

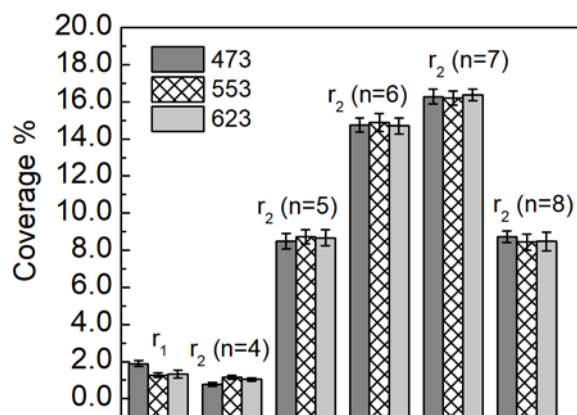


Figure 4.4 Surface coverage for $T_w=473$ K, 553 K and 623 K.

is depicted. 10 s were adequate for the system to reach equilibrium. The growth rate is computed as the average height difference of the surface, $\Delta H_{Cu}/\Delta t$ [102] where the lattice constant between two successive surfaces of Cu is 2.55 Å.

As it can be seen, the predicted growth rate is in satisfactory agreement with the experimental data up until 573 K. Above that T_w the results from the multiscale computations deviate from the experimental values. To better understand the results of Fig. 4.3, we construct the surface coverage i.e., the sites that are available for a process to perform, for 473 K, 553 K and 673 K. The results are shown in Fig. 4.4. When the system reaches equilibrium, in all cases almost half of the surface is idle i.e., half of the sites are blocked because cannot form quartets for the $[Cu(amd)]_2$ molecule to dissociate. The percentages of the idle surface in 473 K, 553 K and 673 K is 49.09 %, 49.33% and 49.41% respectively. Most of the rest of the surface is covered with $Cu_2(amd)$ and amd with coverage values of 49.01 % (473 K), 49.40 % (553 K) and 49.27 % (623 K). The rest - 1.90% (473 K), 1.27 % (553 K) and 1.32 % (623 K) – is available for the dissociation of $[Cu(amd)]_2$. Note that $Cu_2(amd)$ and amd with less than 3 neighbors are instantly desorb. The maximum difference between the surfaces is observed in 473 K where r_1 is ~33% higher compared to the higher temperature cases. Nevertheless, $s_{[Cu(amd)]_2}$ has the lowest value – see Table 4.1 – and for that, the lowest value of the growth is observed. As our model predicts, on the higher temperatures, the increase in T_s is not adequate to increase $s_{[Cu(amd)]_2}$ sufficiently to saturate the surface and the inhibition effects to decrease the deposition rate. For that, most probable cause for the decrease of the growth rate in the high temperatures are the initiation of volumetric reactions in the bulk phase of the CVD reactor with a mechanism as reported in [10]. The initiation of the volumetric reactions causes the decrease of the mole fraction of the $[Cu(amd)]_2$ and consequently to the decrement of the deposition rate.

5. Conclusions

We developed a multiscale computations framework to explore the deposition of Cu from copper(I) N,N'-di-isopropylacetamidinate $[\text{Cu}(\text{iPr-Me-amd})]_2$ or $[\text{Cu}(\text{amd})]_2$ where $\text{amd} = \text{CH}(\text{CH}_3)_2\text{NC}(\text{CH}_3)\text{NCH}(\text{CH}_3)_2$. The multiscale model consists of a macro-scale model which describes the transport phenomena inside the CVD reactor based on the conservation equations (mass, momentum, energy and species mass fractions) and a kMC model which considers inhibition effects from the byproduct, namely H(amd) of the single surface reactor which leads to the deposition of Cu. The link point of the models is the boundary condition for the species equation i.e. the mass fractions and sticking coefficients computed in the macro-scale model is transferred to the kMC model. A single surface reaction is assumed in the macro-scale model according to previous experimental works.

Computations were performed for growth temperatures from 473 K to 623 K. The multiscale model satisfactorily predicts the growth rate of Cu up until 573 K. Above this temperature, the computed growth rate deviates from the experimental data; the multiscale model predicts constant growth rate as the temperature increases while the experimental growth rate decreases. The latter is attributed to the initiation of volumetric reactions above 573 K which can reduce the mole fraction of the precursor leading to the further decrement of the growth rate. Concerning the kMC model, the desorption of H(amd) is assumed to be the limited step and is found to be barrierless (athermal) i.e. it does not depend on temperature and has small activation energy (0.09 eV). The latter, must be validated by DFT computations to validate our kMC model.

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